

PATH INTEGRALS IN QFT

Key object in QFT: evolution operator

$$\langle \psi_f | e^{-iHt} | \psi_i \rangle$$

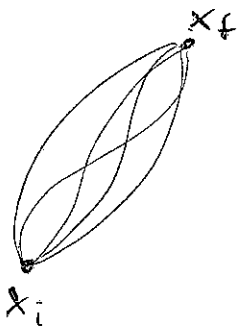
Claim:

$$= \mathcal{N} \int_{\psi(x,t_i) = \psi_i}^{\psi(x,t_f) = \psi_f} [d\psi(x,t)] e^{iS[\psi]/\hbar}$$

↙ action.

e.g. 1 particle in QM

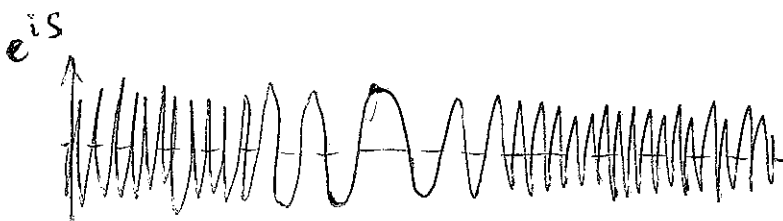
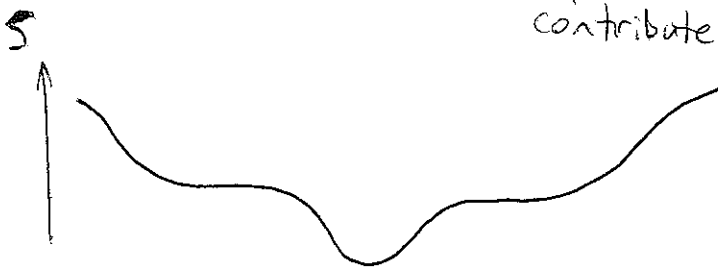
$$\langle x_f | e^{-i(t_f - t_i)H} | x_i \rangle = \mathcal{N} \int_{x(t_i) = x_i}^{x(t_f) = x_f} [dx(t)] e^{iS[x(t)]/\hbar}$$



Quantum amplitude is sum/average over all possible histories weighted by e^{iS}

Classical physics $\Rightarrow \hbar \rightarrow 0$

\therefore integrand rapidly oscillates
- only minimum action trajectories contribute



Derivation: set $\hbar=1$ work with $T \equiv t_f - t_i$

- Break time evolution up into many steps & insert complete basis of states

$$\begin{aligned} \mathcal{U}(x_f, x_i, T) &\equiv \langle x_f | e^{-iHT} | x_i \rangle \\ &= \int \prod dx_n \langle x_f | e^{-iH\epsilon} | x_n \rangle \langle x_n | e^{-iH\epsilon} | x_{n-1} \rangle \dots \langle x_1 | e^{-iH\epsilon} | x_i \rangle \\ &\quad \epsilon \equiv \frac{T}{N} \end{aligned}$$

Q: Calculate $\langle x | e^{-iH\epsilon} | y \rangle$ to $\mathcal{O}(\epsilon)$ for $H = \frac{p^2}{2m} + V(x)$

$$\begin{aligned} &= \int \frac{dp}{(2\pi)} \langle x | e^{-i\hat{H}\epsilon} | p \rangle \langle p | y \rangle \\ &= \int \frac{dp}{(2\pi)} \langle x | e^{-iH(p,x)\epsilon} | p \rangle \langle p | y \rangle \quad \text{ordinary fn.} \\ &= \int \frac{dp}{(2\pi)} e^{-iH(p,x)\epsilon} e^{ip(x-y)} \end{aligned}$$

$$\therefore \mathcal{U} = \int \prod_{n=1}^{N-1} dx_n \prod \frac{dp_n}{(2\pi)} e^{i\epsilon \left(\sum p_j \frac{(x_j - x_{j-1})}{\epsilon} - H(p_j, x_j) \right)}$$

$\epsilon \rightarrow 0$ Define $x(\tau)$ s.t. $x(t_i + \frac{T}{N}n) = x_n$
 $p(\tau)$ s.t. $p(t_i + \frac{T}{N}n) = p_n$

$$\mathcal{U}(x_f, x_i, T) = \int_{x(t_i)=x_i}^{x(t_f)=x_f} \prod_{\tau} [dx(\tau)] \prod_{\tau} \left[\frac{dp(\tau)}{2\pi} \right] e^{i \int_{t_i}^{t_f} dz \{ p(\tau) \dot{x}(\tau) - H(p(\tau), x(\tau)) \}}$$

If $H(x, p) = \frac{p^2}{2m} + V(x)$

get $\int \prod \frac{dp(\tau)}{2\pi} e^{i \int_{t_i}^{t_f} d\tau \left\{ -\frac{1}{2m} (p(\tau) - m\dot{x}(\tau))^2 \right\}} e^{i \int_{t_i}^{t_f} d\tau \left\{ \frac{m}{2} \dot{x}^2(\tau) - V(x) \right\}}$

$\therefore \mathcal{U} = \mathcal{N} \int_{x(t_i)=x_i}^{x(t_f)=x_f} \prod [dx(\tau)] e^{i \int_{t_i}^{t_f} d\tau \left[\frac{m}{2} \dot{x}^2(\tau) - V(x(\tau)) \right]}$
 e^{iS}

Correlation functions:

$\langle x_f, t_f | \prod \{ \theta_n(x(t_n)) \dots \theta_1(x(t_1)) \} | x_i, t_i \rangle$
 $= \langle x_f, t_f | \theta_n(x(t_n)) \dots \theta_1(x(t_1)) | x_i, t_i \rangle$ (assume $t_f \geq t_n \geq t_{n-1} \dots \geq t_i$)
 $= \mathcal{N} \int_{x(t_i)=x_i}^{x(t_f)=x_f} \prod [dx(\tau)] \theta_n(x(t_n)) \dots \theta_1(x(t_1)) e^{iS}$

↑ ↑
ordinary fns:
order doesn't matter.

Path integral averages automatically give time-ordered correlation fns.

Field theory: want

$\langle \Omega | T \{ \theta_1(x(t_1)) \dots \theta_n(x(t_n)) \} | \Omega \rangle$ ← vacuum.

Peskin trick $e^{-iHT} |\psi\rangle = \sum_n e^{-iE_n T} |n\rangle \langle n | \psi \rangle$

$T \rightarrow \infty(1-i\epsilon)$ picks out term w. lowest energy

$\therefore \lim_{T \rightarrow \infty(1-i\epsilon)} e^{-iHT} |\psi\rangle = c |\Omega\rangle$

$$\langle \Omega | T \{ \theta_1 \dots \theta_n \} | \Omega \rangle$$

$$= \lim_{T \rightarrow \infty (1-i\epsilon)} \frac{\int [d\phi_i(\tau)] e^{i \int_{-T}^T L[\phi_i(\tau)] d\tau} \theta_1(\phi(\tau_1)) \dots \theta_n(\phi(\tau_n))}{\int [d\phi(\tau)] e^{i \int_{-T}^T L[\phi(\tau)] d\tau}}$$

↑
cancels out normalization factor.

How do we calculate this?

$$\mathcal{L} = \mathcal{L}_0 + \mathcal{L}_{int}$$

↑
QUADRATIC

$$\langle \Omega | T \{ \theta_1 \dots \theta_n \} | \Omega \rangle = \lim \frac{\int [dx(\tau)] e^{i \int \mathcal{L}_0 dx} \theta_1 \dots \theta_n \left(\sum \frac{i^n}{n!} (\int \mathcal{L}_{int})^n \right)}{\int [dx(\tau)] e^{i \int \mathcal{L}_0[x(\tau)]} \left(\sum \frac{i^n}{n!} (\int \mathcal{L}_{int})^n \right)}$$

sum of integrals like

$$\int_{-\infty}^{\infty} dx_i e^{-\frac{1}{2} M_{ij} x_i x_j} x_i^{i_1} \dots x_i^{i_n}$$