

LAST TIME: UV divergences in higher-order amplitudes

solution: include cutoff

$$X + (\cancel{X} + \cancel{X} + \dots)$$

$$S(m, \lambda) \rightarrow S(m, \lambda, \Lambda)$$

$$\phi(x) \rightarrow \phi_\Lambda(x) = \int_{|\vec{p}| < \Lambda} d^3\vec{p} e^{i\vec{p}\cdot\vec{x}} \phi(\vec{p})$$

can define physical parameters  $m_{\text{PHYS}}, \lambda_{\text{PHYS}}(m, \lambda, \Lambda)$  so that any low-energy observable

$$\sigma(m, \lambda, \Lambda) = \sigma(m_{\text{PHYS}}, \lambda_{\text{PHYS}}) + \mathcal{O}\left(\frac{1}{\Lambda}\right)$$

↑ can take  $\Lambda \rightarrow \infty$  to get finite result.

$m_{\text{PHYS}}$  = physical mass  
= pole in full propagator.

$$\begin{array}{c} \text{leading} \\ \text{order} \end{array} \quad \frac{1}{p^2 - m^2} + \dots$$

↑  
pole at  $p^2 = m^2$

full result: pole at  $\frac{1}{p^2 - m^2} \rightarrow p^2 = m_{\text{PHYS}}^2 = m^2 + \lambda \Lambda^2 + \dots$

$\lambda_{\text{PHYS}}$ : full scattering amplitude at  $p \rightarrow 0$

$$X + (\cancel{X} + \dots) + \dots$$

$$\lambda_{\text{PHYS}} = \lambda + \lambda^2 \ln(\Lambda) + \dots$$

# MORE GENERAL:

- Can start with

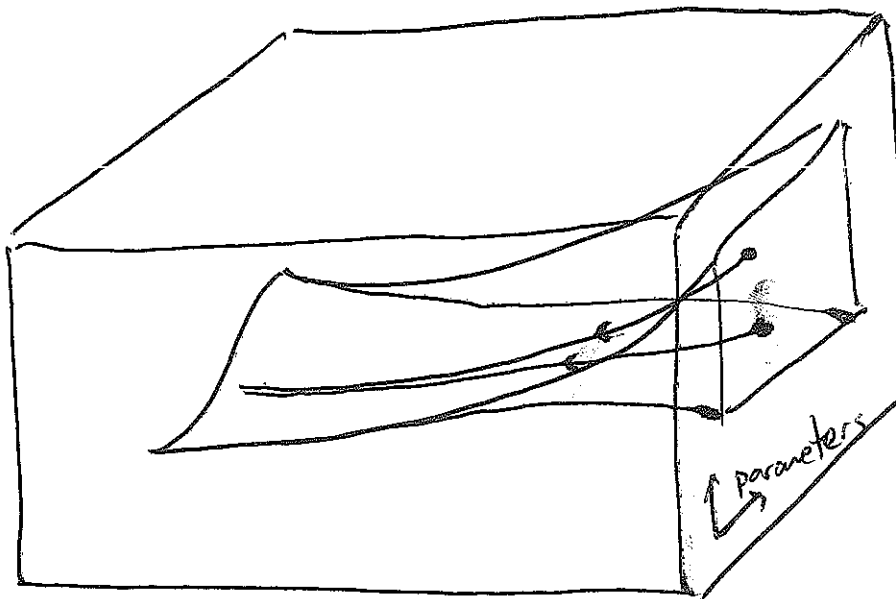
$$S[m_i, \lambda_i, \phi_\Lambda]$$

↑ all possible couplings

- low-energy physics  $E \ll \Lambda$  depends only on small number of physical parameters

$$\sigma(m_i, \lambda_i, \Lambda) = \underbrace{\sigma(m_{\text{phys}}^i, \lambda_{\text{phys}}^a)}_{\text{small \# parameters}} + \mathcal{O}\left(\frac{1}{\Lambda}\right)$$

large # parameters



- many UV theories have same IR physics.

- can find simple theory with cutoff  $\Lambda \gg E$  that has same low E predictions as real theory.

→ cutoff  $\Lambda$

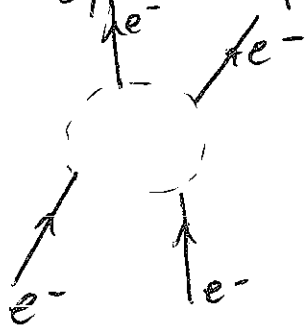
# IR DIVERGENCES:



$\int d^4k$ . divergent even with cutoff  $|k| < \Lambda$

→ These cancel when we include all diagrams relevant to a physical measurement.

Real experiment: particle detectors have lower limit on energy of photons detected



indistinguishable from

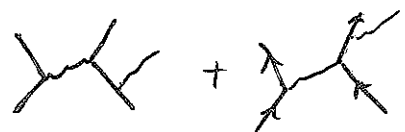
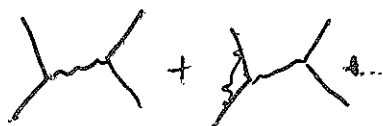


if  $E_{\text{photon}} < E_{\text{min}}$

Really want cross section for

$e^-e^- \rightarrow e^-e^- + \text{arbitrary number of "soft" photons}$   
 $E < E_{\text{min}}$

$$\sigma = \sigma(e^-e^- \rightarrow e^-e^-) + \int_0^{E_{\text{min}}} dE \sigma(e^-e^- \rightarrow e^-e^- \gamma(E))$$



← divergences cancel. →

Full answer for  $e^-$  scattering off of very heavy charged particle:

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{meas}} = \left(\frac{d\sigma}{d\Omega}\right)_0 \times \exp\left(-\frac{\alpha}{\pi} f(\vec{p}-\vec{q}) \ln\left(\frac{(\vec{p}-\vec{q})^2}{E_{\text{min}}}\right)\right)$$

$$\rightarrow 0 \quad \text{when } E_{\text{min}} \rightarrow 0$$

Cross section with no photons = 0

ACCELERATED CHARGES RADIATE!

radiated photons = BREHMSSTRAHLUNG.