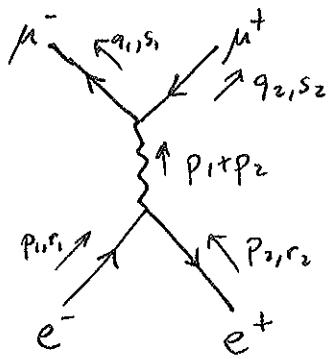


LAST TIME: cross section for $e^+e^- \rightarrow \mu^+\mu^-$



$$iM = \frac{ie^2}{(p_1+p_2)^2} \bar{u}^{s_1}(q_1) \gamma^\mu v^{s_2}(q_2) \bar{v}^{r_2}(p_2) \gamma_\mu u^{r_1}(p_1)$$

for cross section need $|M|^2 = M M^*$

AVERAGE over initial spins (if don't know)

SUM over final spins (if don't measure)

$$\text{NEED: } (\bar{\psi} \gamma^\mu \psi)^* = \bar{\psi} \gamma^\mu \gamma^\nu \gamma^\lambda \psi$$

$$= \bar{\psi} \gamma^\lambda \gamma^\nu \gamma^\mu \psi$$

$$\Rightarrow (\bar{\psi} \gamma^\mu \psi)^* = \bar{\psi} \gamma^\mu \psi$$

$$\text{Want: } \frac{1}{2} \sum_{r_1} \frac{1}{2} \sum_{r_2} \sum_{s_1, s_2} |M_{r_1, r_2, s_1, s_2}|^2$$

$$= \frac{1}{4} \sum_{\substack{r_1, r_2 \\ s_1, s_2}} \bar{u}^{s_1}(q_1) \gamma^\mu v^{s_2}(q_2) \bar{v}^{s_2}(q_2) \gamma^\nu u^{s_1}(q_1) \\ \bar{v}^{r_2}(p_2) \gamma_\nu u^{r_1}(p_1) \bar{u}^{r_1}(p_1) \gamma_\nu v^{r_2}(p_2)$$

Spin sums: use:

$$\sum_s u_\alpha^s(p) \bar{u}_\beta^s(p) = (\not{p} + m)_{\alpha\beta}$$

$$\sum_s v_\alpha^s(p) \bar{v}_\beta^s(p) = (\not{p} - m)_{\alpha\beta}$$

$$\text{Get e.g. } \bar{u}_\alpha^{s_1}(q_1) \gamma^\mu_{\alpha\beta} v_\beta^{s_2}(q_2) \bar{v}_\rho^{s_2}(q_2) \gamma^\nu_{\rho\sigma} u_\sigma^{s_1}(q_1)$$

$$\rightarrow \gamma^\mu_{\alpha\beta} (\not{q}_2 - m)_{\beta\rho} \gamma^\nu_{\rho\sigma} (\not{q}_1 + m)_{\sigma\alpha}$$

$$= \text{Tr}(\gamma^\mu (\not{q}_2 - m) \gamma^\nu (\not{q}_1 + m))$$

$$= 4 (q_1^\mu q_2^\nu + q_1^\nu q_2^\mu - q_1 \cdot q_2 \eta^{\mu\nu} - m_\mu^2 m_\nu^{\mu\nu})$$

Need:

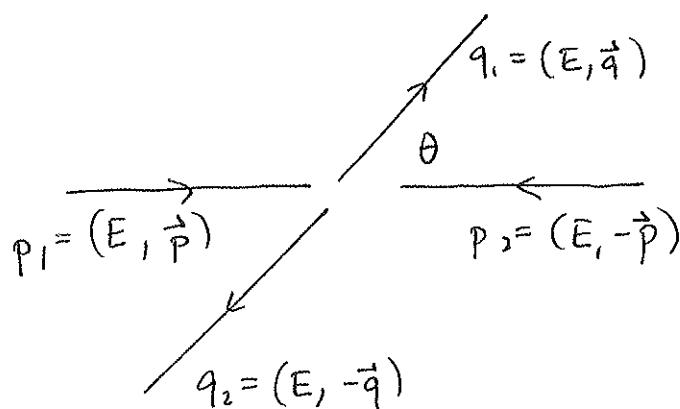
$$\text{Tr}(\gamma^\mu) = 0 \quad \text{Tr}(\gamma^\mu \gamma^\nu) = 4 \eta^{\mu\nu} \quad \text{Tr}(\gamma^\mu \gamma^\nu \gamma^\lambda) = 0$$

$$\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\lambda \gamma^\sigma) = 4 (\eta^{\mu\nu} \eta^{\lambda\sigma} + \eta^{\nu\lambda} \eta^{\sigma\mu} - \eta^{\mu\lambda} \eta^{\nu\sigma})$$

$$\frac{1}{4} \sum |M|^2 = \frac{8e^4}{(p_1 + p_2)^4} [(p_1 \cdot q_1)(p_2 \cdot q_2) + (p_1 \cdot q_2)(p_2 \cdot q_1) + m_\mu^2 (p_1 \cdot p_2) + m_e^2 (q_1 \cdot q_2)]$$

choose frame: C.o.m. simplest

must have $E_e \gg m_e$ so approximate $m_e = 0$
 $\therefore E \approx |\vec{p}_1| = |\vec{p}_2|$



$$\sqrt{|M|^2} = \frac{8e^4}{(2E)^4} \left((E^2 - E|\vec{q}| \cos\theta)^2 + (E^2 + E|\vec{q}| \cos\theta)^2 + 2m_\mu^2 E^2 \right)$$

$$= e^4 \left[\left(1 + \frac{m_\mu^2}{E^2} \right) + \left(1 - \frac{m_\mu^2}{E^2} \cos^2\theta \right) \right]$$

Plug into cross section formula, integrate over all energies/momenta that are fixed by 8 fns:

$$d\sigma = \frac{1}{2E_{p_1} \cdot 2E_{p_2} \cdot |v_1 - v_2|} \left(\frac{1}{4} \sum |M|^2 \right) \underbrace{\frac{d^3 q_1}{(2\pi)^3 2E_{q_1}} \frac{d^3 q_2}{(2\pi)^3 2E_{q_2}} (2\pi)^4 \delta^4(p_1 + p_2 - q_1 - q_2)}_{\text{same as for } d\Phi \rightarrow \Phi\Phi}$$

$$\text{gives } \frac{191}{(2\pi)^2 \cdot 4 \cdot E_{cm}} \cdot d\Omega$$

$$\frac{d\sigma}{d\Omega} = \left(\frac{e^2}{4\pi} \right)^2 \frac{1}{16E^2} \sqrt{1 - \frac{m_\mu^2}{E^2}} \left[\left(1 + \frac{m_\mu^2}{E^2} \right) + \left(1 - \frac{m_\mu^2}{E^2} \right) \cos^2\theta \right]$$

$$\boxed{\sigma_{tot} = \int_{\text{sphere}} \frac{d\sigma}{d\Omega} \cdot d\Omega = \frac{4\pi \alpha^2}{3 E_{cm}} \sqrt{1 - \frac{m_\mu^2}{E^2}} \left(1 + \frac{1}{2} \frac{m_\mu^2}{E^2} \right)}$$

→ excellent agreement w. experiment

→ can be used to measure mass of particle.

~~Q2 How does p₁ relate to p₂ in 2-body~~

very large E: $\sigma \sim \frac{1}{E^2}$ (dimensional analysis)
 if m_e, m_μ irrelevant.

$\frac{1}{M}$ \leftarrow width
 $E \rightarrow p$ \leftarrow particle mass

"Resonance": signature of new unstable particle



$$W \gg J \quad \text{if} \quad \frac{\frac{J}{M^2} + (E-M)^2}{1} \sim 0$$

\downarrow
 decay to other
 particles

$$e^{-imt} \xrightarrow{\text{evolution}} e^{-imt - \frac{J^2}{M^2}}$$

shifts $M \rightarrow M - \frac{J^2}{2}$ in propagator

$$\xrightarrow{\text{decay rate}} \dots + \text{---} + \dots$$

since particle unstable.

Higher orders smooths this out



③ $E_m = M$ result follows up

② $E_m \ll M$ mass unimportant

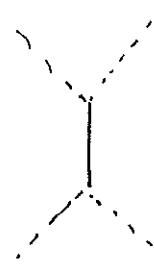
physics same if we replace $m\phi^2 \rightarrow \frac{M^2}{\phi^2}$ with $\frac{M^2}{\phi^2}$

① $E_m \gg M$ propagator acts like const. $-\frac{1}{M^2}$

final result has factor $\frac{(E_m - M^2)}{1}$

$$\frac{1}{(p_1 + p_2)^2 - M^2} \xleftarrow{\text{propagator}} \frac{(p_1 + p_2)}{1}$$

e.g. HW



Interactions via massive particle