

## Problem set 2

1. Draw all distinct proper self-energy diagrams of 3rd order, and find their multiplicities, for disorder averaged Green's functions.

2. Use path integral techniques to find the propagator  $\mathcal{G}(xt, x_0t_0)$  for a charged particle moving in a uniform electric field, described by  $\hat{H} = \frac{\hat{p}^2}{2m} - qE\hat{x}$ . Are any approximations needed?

3. In class, we derived the propagator for a free particle to be

$$\mathcal{G}_0(xt, x_0t_0) = \sqrt{\frac{m}{2\pi i\hbar\Delta t}} \exp\left(\frac{im}{2\hbar\Delta t}(x - x_0)^2\right)$$

where  $\Delta t = t - t_0$ . Consider now a free particle moving on a finite chain of length  $L$ , which has periodic boundary conditions (in other words, a circle).

(a) Find the exact eigenspectrum and eigenvalues and derive an expression for the propagator starting from the definition

$$\mathcal{G}(xt, x_0t_0) = \langle x | \exp\left(-\frac{i}{\hbar}\hat{H}(t - t_0)\right) | x_0 \rangle$$

(b) use path integral techniques to derive the propagator. Hint: in a periodic space there may be more than one way to get from "here" to "there".

You might find this useful: the Jacobi theta function is defined by

$$\theta_3(z, t) = \sum_{n=-\infty}^{\infty} \exp\left(i\pi t n^2 + i2nz\right)$$

and satisfies the properties

$$\begin{aligned} \theta_3(z, t) &= \theta_3(z + \pi, t) \\ \theta_3(z + \pi t, t) &= e^{-i\pi t - 2iz} \theta_3(z, t) \\ \theta_3(z, t) &= (-it)^{-\frac{1}{2}} e^{z^2/i\pi t} \theta_3\left(\frac{z}{t}, -\frac{1}{t}\right) \end{aligned}$$

4. Find the ground-state wavefunction for a harmonic oscillator  $\langle x | n = 0 \rangle = \phi_0(x)$  starting from  $a|0\rangle = 0$ . Write  $a$  in terms of  $\hat{x}$  and  $\hat{p}_x$ , and project the equation on  $\langle x |$  to get a differential equation for  $\phi_0(x)$ . Solve it and normalize the solution. Note: you can now use the same method starting from  $a^\dagger|n - 1\rangle = \sqrt{n}|n\rangle$  to find all eigenstates  $\langle x | n \rangle$ .

5. Use the same approach as in (4) to find (up to an overall phase shift)  $\langle x | z \rangle$ , starting from the fact that  $a|z\rangle = z|z\rangle$ . Check your result by computing this wavefunction directly from definition

$$\langle x | z \rangle = e^{-\frac{|z|^2}{2}} \langle x | e^{\frac{z}{\sqrt{2i}}(\hat{x} - i\hat{p}/(m\omega))} | 0 \rangle$$

Remember that  $T(a) = \exp(ia\hat{p}/\hbar)$  is simply the translation operator:  $\langle x | T(a) | \psi \rangle = \psi(x + a)$  (you should be able to check this). Using this result and  $\exp(-\frac{i}{\hbar}\hat{H}t)|z\rangle = \exp(-i\omega t/2)|z\rangle \exp(-i\omega t)$ , check the equations written directly above Eq. (1.64) in the notes.