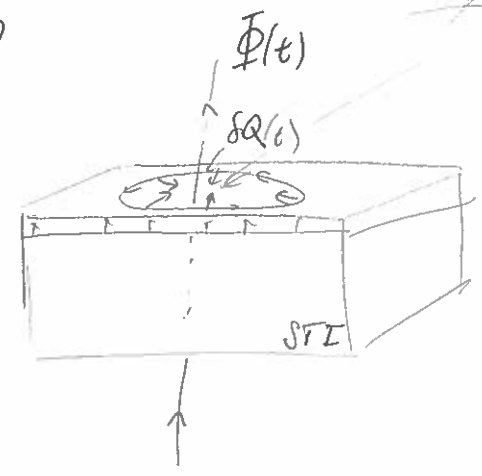


Flux tube insertion through experiment, the "wormhole effect"

[Rosenberg, Guo, Franz, PRB 82, 041104R (2010)]

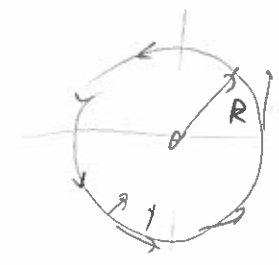
Faraday's law:

$$\vec{\nabla}_x \cdot \vec{E} = -\frac{1}{c} \frac{\partial B}{\partial t} \quad \int dS$$



Hall current response

$$\vec{j} = \sigma_{xy} (\vec{E} \times \hat{z})$$



$$\oint \vec{E} \cdot d\vec{l} = -\frac{1}{c} \frac{\partial \Phi}{\partial t} \Rightarrow$$

$$= 2\pi R E = -\frac{1}{c} \frac{\partial \Phi}{\partial t}$$

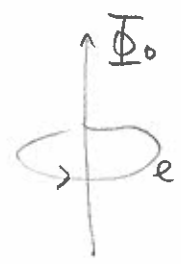
$$j = \sigma_{xy} \frac{1}{2\pi R} \left(-\frac{1}{c} \frac{\partial \Phi}{\partial t} \right)$$

$$\sigma_{xy} = \frac{e^2}{2h}, \quad \Phi_0 = \frac{hc}{e}$$

$$\delta Q = \int_0^T dt \quad 2\pi R j = \sigma_{xy} \frac{\delta \Phi}{c}$$

Choose $\delta \Phi = \Phi_0$

$$\delta Q = \frac{e^2}{2h} \frac{hc}{e} = \frac{1}{2} e$$



Phase 2π
 $e^{i2\pi} = 1$

However, a flux tube with Φ_0 flux can be removed from the Hamiltonian by a gauge twist.

This leads to a paradox:

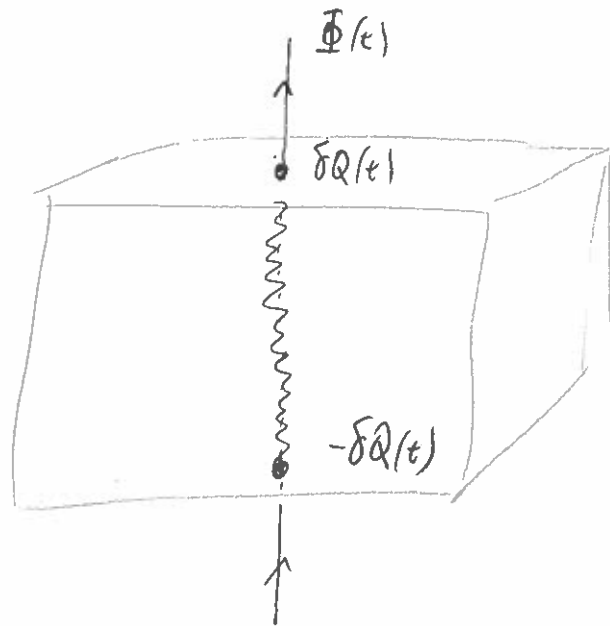
• a state with charge $\pm \frac{1}{2}e$ should be an eigenstate of a non-interacting Hamiltonian.

Solution:

"Wormhole effect"

When $\Phi(t) = \frac{1}{2}\Phi_0$ there is a gapless mode formed along the tube inside STI.

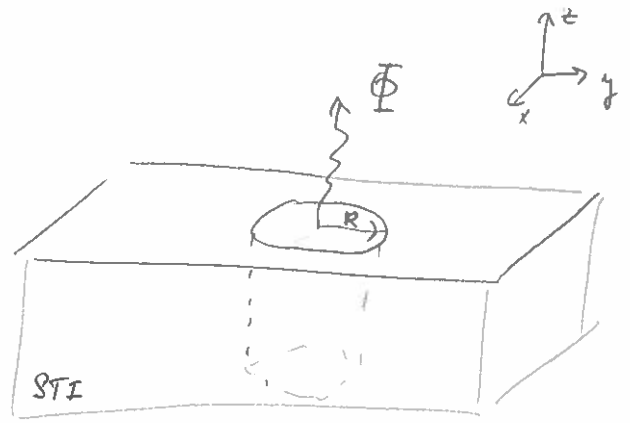
The charge $\delta Q(t)$ escapes along the wormhole to the other side.



• Image monopole effect [Qi, Li, Zhang & Zhang, *Science* 323, 1184 (2005)]

Wormhole effect calculation

According to Ostrowsky, Boruyi & Hirlein
[REL 105, 036803 (2010), Supplement]
surface states of a STI in
a curved geometry are
governed by



$$\mathcal{H} = \frac{i}{2} [\vec{S} \cdot \vec{p}, \vec{S} \cdot \hat{n}] = \frac{\vec{\nabla} \cdot \hat{n}}{2} + \frac{1}{2} \{ \hat{n} \cdot (\vec{p} \times \vec{S}) + (\vec{p} \times \vec{S}) \cdot \hat{n} \}$$

where $\vec{p} = -i\vec{\nabla}$ and \hat{n} is a surface normal unit vector.

• For cylindrical hole

$$\hat{n} = - \left(\frac{x}{R}, \frac{y}{R}, 0 \right)$$

$$\cdot (\vec{p} \times \vec{S}) \cdot \hat{n} = - (\vec{S} \times \vec{p}) \cdot \hat{n} = - \vec{S} \cdot (\vec{p} \times \hat{n})$$

$$\vec{\nabla} \times \hat{n} = 0, \quad \vec{\nabla} \cdot \hat{n} = - \frac{1}{R}$$

$$\mathcal{H} = - \frac{1}{2R} + \vec{S} \cdot (\hat{n} \times \vec{p}) \quad \text{For cylindrical hole}$$

Minimal substitution to include mag. flux

$$\vec{p} \rightarrow \vec{\pi} = \vec{p} - \frac{e}{c} \vec{A}$$

$$\vec{A} = \frac{\hat{z} \times \vec{r}}{2\pi r} \Phi_0 = (-y, x, 0) \frac{\Phi_0}{2\pi r^2}$$

Φ = Φ₀ in units of Φ₀ = h/e²

Transform to cylindrical coordinates

$$x = r \cos \varphi$$

$$y = r \sin \varphi$$

$$z = z$$

$$\hat{n} = -(\cos \varphi, \sin \varphi, 0)$$

$$\hat{n} \times \hat{\Pi} = - \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \cos \varphi & \sin \varphi & 0 \\ \pi_x & \pi_y & \pi_z \end{vmatrix} = (-\hat{x} \sin \varphi + \hat{y} \cos \varphi) \pi_z - \hat{z} (\cos \varphi \pi_y - \sin \varphi \pi_x)$$

$$\left[\mathcal{H} = -\frac{1}{2R} + (-s_x \sin \varphi + s_y \cos \varphi) (-i\partial_z) - s_z [\cos \varphi (-i\partial_y - A_y) - \sin \varphi (-i\partial_x - A_x)] \right]$$

$$\partial_x = \cos \varphi \partial_r - \sin \varphi \frac{1}{r} \partial_\varphi$$

$$\partial_y = \sin \varphi \partial_r + \cos \varphi \frac{1}{r} \partial_\varphi$$

$$\frac{e}{c} A_x = -\frac{\sin \varphi}{r} z \quad \frac{e}{c} A_y = \frac{\cos \varphi}{r} z$$

• Assume plane wave solutions along z : $-i\partial_z \rightarrow k$

$$\mathcal{H}_k = -\frac{1}{2R} + \begin{pmatrix} \frac{1}{R} (i\partial_\varphi + k) & -ik e^{-i\varphi} \\ ik e^{i\varphi} & -\frac{1}{2} (i\partial_\varphi + k) \end{pmatrix}$$

$$\mathcal{H}_k \Psi_k(\varphi) = E_k \Psi_k(\varphi)$$

Assume solution

$$\Psi_k(\varphi) = e^{i\varphi l} \begin{pmatrix} f_k \\ e^{i\varphi} g_k \end{pmatrix} \quad l = 0, \pm 1, \pm 2, \dots$$

$$= e^{i\varphi l} \begin{pmatrix} 1 & 0 \\ 0 & e^{i\varphi} \end{pmatrix} \begin{pmatrix} f_k \\ g_k \end{pmatrix}$$

$$\tilde{\mathcal{H}}_k \tilde{\Psi}_k = E_k \tilde{\Psi}_k \quad \tilde{\Psi}_k = \begin{pmatrix} f_k \\ g_k \end{pmatrix}$$

$$\tilde{\mathcal{H}}_k = s_y k - s_z \frac{1}{R} \left(l + \frac{1}{2} - \nu \right)$$

$$l = 0, \pm 1, \pm 2 \dots$$

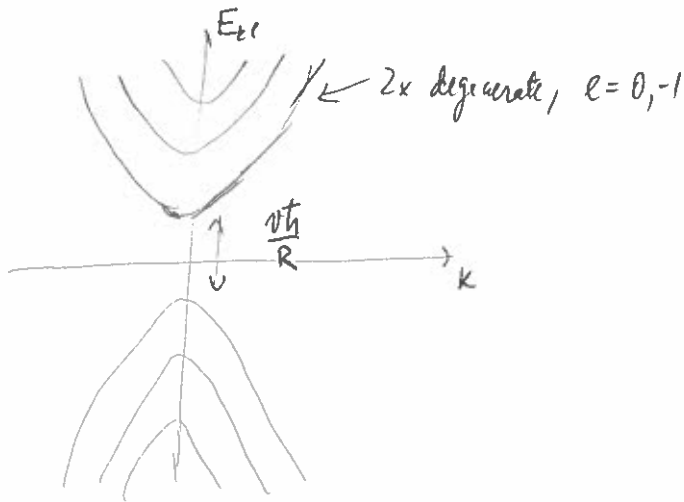
$$\nu = \frac{\Phi}{\Phi_0}$$

$$\tilde{\mathcal{H}}_{k\nu} = U_\nu^{-1} \mathcal{H}_k U_\nu$$

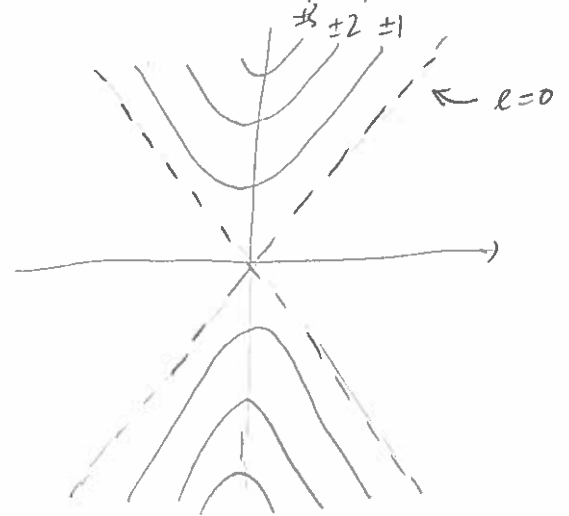
$$E_{k\nu} = \pm v \hbar \sqrt{k^2 + \frac{1}{R^2} \left(l + \frac{1}{2} - \nu \right)^2}$$

$$l = 0, \pm 1, \pm 2 \dots$$

$\nu = 0$ (no flux)



$\nu = \frac{1}{2}$ (half-flux g.)



$$\nu = \frac{1}{2} : E_{k0} = \pm v \hbar |k| \quad \text{indep. of } R$$

↓ "Wormhole effect"

Back to the Flux insertion exp.:

