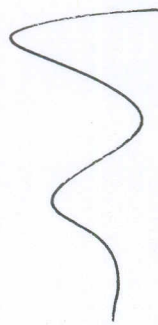


» PHYS 501 «

Solutions

of

HW # 4



Q1.

$$\vec{S} = \hbar \vec{\sigma}, \quad |\vec{\sigma}| = \frac{1}{2}$$

$$H = -\vec{\mu} \cdot \vec{B}, \quad \vec{\mu} = g\mu_B \vec{S} = g\mu_B \hbar \vec{\sigma}$$

$$\vec{B} = \hat{z} B_0 + \hat{x} b_1 f(t)$$

$$\Rightarrow H = \underbrace{-g\mu_B \hbar B_0 \sigma_z}_{\hat{H}_0} - \underbrace{b_1 g\mu_B \hbar f(t) \sigma_x}_{\hat{V}(t)}$$

$$C_n(t) = C_n(t \rightarrow -\infty) - \frac{i}{\hbar} \int_{-\infty}^t (V_I(t'))_{nm} C_m(t') dt'$$

$nm \rightarrow$ summation

first order:

$$C_n(\infty) = C_n(-\infty) - \frac{i}{\hbar} \int_{-\infty}^{\infty} (V_I(t'))_{nm} C_m(-\infty) dt'$$

two-level $\Rightarrow n=1,2 \quad \hat{z}\uparrow \rightarrow 1, \hat{z}\downarrow \rightarrow 2$

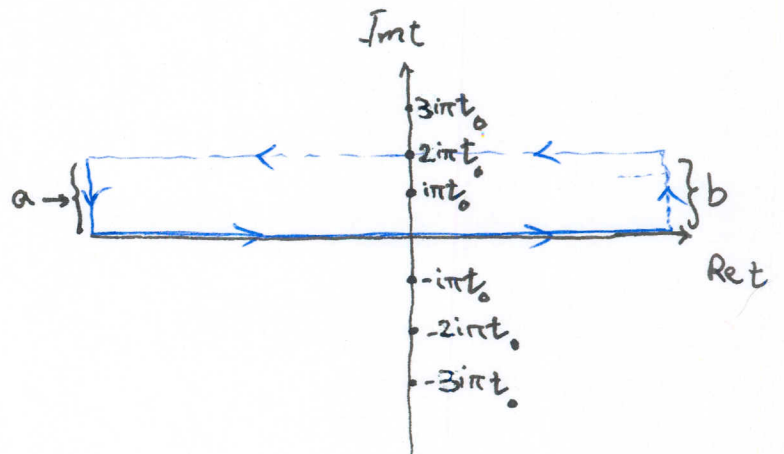
$$C_n(-\infty) = \delta_{n1}$$

$$\Rightarrow C_2(\infty) = \cancel{\delta_{21}} - \frac{i}{\hbar} \int_{-\infty}^{\infty} (V_I(t'))_{21} dt'$$

$$(V_I(t'))_{21} = e^{ig\mu_B B_0 t'} \cdot \underbrace{(-b_1 g\mu_B \hbar f(t))}_C = C e^{i\omega t'} f(t')$$

$$\Rightarrow C_2(\omega) = -\frac{i}{h} C \int_{-\infty}^{\infty} e^{i\omega t'} f(t') dt'$$

$$f(t') = \frac{e^{t'/t_0}}{1 + e^{t'/t_0}}$$



$$\oint f(z) dz = 2\pi i \text{Res}(i\pi)$$

Since $f(z)$ on the a or b sides have nonzero value we should be careful when we integrate,

Note that $f(t) = \frac{1}{1 + e^{-t/t_0}} = 1 - \frac{1}{1 + e^{t/t_0}} = 1 - f(-t)$

$$\Rightarrow f(t) + f(-t) = 1 \xrightarrow{\text{F.T}} f(\omega) + f^*(\omega) = 2\pi \delta(\omega)$$

$\Rightarrow \text{Re}f(\omega) = \pi \delta(\omega) \Rightarrow$ This in infact comes from the a, b side integration

for the two other sides we have:

$$P \Rightarrow \left[\int_{-\infty}^{\infty} e^{i\omega t} f(t) dt - e^{-2\pi i \omega t_0} \int_{-\infty}^{\infty} e^{i\omega t} f(t) dt \right] = 2\pi i \cdot (i e)^{\cdot} \cdot (-t_0)$$

excluding the Dirac delta function contribution

$$\Rightarrow \sinh \omega \pi t_0 \int_{-\infty}^{\infty} e^{i\omega t} f(t) dt = -\pi t_0$$

$$\Rightarrow \int_{-\infty}^{\infty} e^{i\omega t} f(t) dt = \frac{i\pi t_0}{\sinh \omega \pi t_0} + \pi \delta(\omega)$$

$$\Rightarrow C_2(\omega) = -\frac{i}{\hbar} C \cdot \left[\frac{i\pi t_0}{\sinh[\omega \pi t_0]} + \pi \delta(\omega) \right]$$

$$\omega = g \mu_B |B_0|, \quad C = -b_1 g \mu_B \hbar$$

$$\Rightarrow C_2(\omega) = \frac{b_1 g \mu_B \pi t_0}{\sinh[g \mu_B |B_0| \pi t_0]}$$

(ii) if $g \mu_B |B_0| t_0 \ll 1$ it means that $\omega t_0 \ll 1$

This means that for our system $f(t)$ behaves like a step function $\Theta(t)$ so we have

$$C_2(t) = -\frac{i}{\hbar} C \int_{-\infty}^t \Theta(t') e^{i\omega t'} dt' = -\frac{i}{\hbar} C \left[\frac{e^{i\omega t} - 1}{i\omega} \right] \Theta(t)$$

$$\Rightarrow C_2(t) = \frac{b_1 g k_B}{g k_B |B_0|} \begin{bmatrix} e^{i g k_B |B_0| t} & \\ & -1 \end{bmatrix} \Theta(t)$$

So the system remains into initial state for $t < 0$ and from $t = 0$ it starts + transition to \uparrow state with the amplitude given above.

$$H_0 = -K_0 \hat{S}_z^2 + E_0 \hat{S}_x^2$$

$$K_0, E_0 > 0$$

$$(i) E_0 = 0 \Rightarrow H = -K_0 \hat{S}_z^2$$

$$\hat{S}_z |M\rangle = \hbar M |M\rangle$$

$$M = -S, -S+1, \dots, S-1, S$$

$$\Rightarrow E_M = -K_0 \hbar^2 M^2, \text{ so } |M\rangle \& | -M\rangle \text{ are degenerate.}$$

$$(ii) \frac{E_0}{K_0} \ll 1, \quad \hat{S}_x^2 = \left(\frac{S_+ + S_-}{2} \right)^2 = \frac{S_+^2 + S_-^2 + S_+ S_- + S_- S_+}{4}$$

$$\Delta_M^S = 2 \cdot \left\{ V^{-M, -(M-2)} \cdot \frac{1}{|E_M^0 - E_{-(M-2)}^0|} \cdot V^{-(M-2), (M-4)} \cdot \frac{1}{|E_{M-2}^0 - E_{-(M-4)}^0|} \cdots V^{M-4, M-2} \cdot \frac{1}{|E_M^0 - E_{-(M-2)}^0|} \cdot V^{M-2, M} \right\}$$

$$\Delta E_M(k) \equiv |E_M^0 - E_{-(M-2k)}^0| \quad k=1, \dots, M-1$$

$$V_M(k) = V^{-(M-2k+2), -(M-2k)} = \langle -(M-2k+2) | (+E_0 \hat{S}_x^2) | -(M-2k) \rangle$$

$$\Rightarrow \Delta E_M(k) = +K_0 \hbar^2 [M^2 - (M-2k)^2] = +4K_0 \hbar^2 k(M-k)$$

$$V_M(k) = E_0 \hbar^2 \left[(S+M-2k+2)(S+M-2k+1)(S-M+2k)(S-M+2k-1) \right]^{\frac{1}{2}}$$

$$\Rightarrow \Delta_M^S = 2 \cdot \left[\prod_{k=1}^{M-1} \frac{V_M(k)}{\Delta E_M(k)} \right] \cdot V^{-M, -(M-2)}$$

$$\Rightarrow \Delta_m^- = \frac{2 \cdot V}{(4K_0)^{M-1}} \cdot \left(\prod_{k=1}^{M-1} \frac{1}{k(M-k)} \right) \cdot \left(\frac{E_0}{4} \right) \cdot \left(\prod_{k=1}^{M-1} \tilde{V}_m(k) \right)$$

$$\downarrow$$

$$\frac{1}{[(M-1)!]^2}$$

where $\tilde{V}_m(k) = \frac{4V_m(k)}{E_0 \hbar^2}$

$$V \prod_{k=1}^{M-1} \tilde{V}_m(k) = \prod_{k=1}^M \tilde{V}_m(k) = \frac{(S+M)!}{(S-M)!}$$

$$\Rightarrow \Delta_m^S = 2 \cdot \left(\frac{E_0}{16K_0} \right)^{M-1} \cdot \frac{E_0}{4} \cdot \frac{(S+M)!}{[(M-1)!]^2 (S-M)!}$$

$$\Rightarrow \Delta_m^S = A_m^S K_0 \left(\frac{E_0}{16K_0} \right)^M$$

$$A_m^S = \frac{8}{[(M-1)!]^2} \cdot \frac{(S+M)!}{(S-M)!}$$

Stirling's approx:

$$\lim_{n \rightarrow \infty} \left[\frac{n!}{\sqrt{2\pi n} \left(\frac{n}{e} \right)^n} \right] = 1$$

$$A_m^S = \left[\frac{8}{((S-1)!)^2} \cdot (2S)! \right] = \frac{8 \cdot \sqrt{2\pi \times (2S)} \cdot \left(\frac{2S}{e} \right)^{2S}}{\left[\sqrt{2\pi(S-1)} \cdot \left(\frac{S-1}{e} \right)^{S-1} \right]^2}$$

$$\Rightarrow A_s^S = \frac{8 \cdot \sqrt{4\pi S} \cdot 4 \cdot S^{2S} \cdot e^{-2S}}{[\sqrt{2\pi S} \cdot S^{S-1} \cdot e^{-S+1}]^2} = \frac{2 \cdot 4^{S+1}}{\sqrt{\pi S}} \cdot \left(\frac{S}{e}\right)^2$$

$$\Rightarrow \Delta_s^S = \frac{2 \cdot 4^{S+1} \cdot K_0 \cdot S^{3/2}}{\sqrt{\pi} \cdot e^2} \cdot \left(\frac{E_0}{16K_0}\right)^S$$

$$\ln \left[\frac{\Delta_s^S}{K_0} \right] = \ln C + \frac{3}{2} \ln S + S \ln \left(\frac{E_0}{4K_0} \right)$$

where $C = \frac{8}{\sqrt{\pi}} \cdot \frac{1}{e^2}$