

#1. (i)

$$\begin{aligned}\langle X_N, t_N | X_1, t_1 \rangle &= \int dX_{N-1} \langle X_N, t_N | X_{N-1}, t_{N-1} \rangle \langle X_{N-1}, t_{N-1} | X_1, t_1 \rangle \\ &= \int dX_{N-1} \sqrt{\frac{m}{2\pi i \hbar \Delta t}} \exp \left[ \left( \frac{i m}{2\hbar} \right) \frac{(X_N - X_{N-1})^2}{\Delta t} + \frac{i}{\hbar} q A (X_N - X_{N-1}) \right. \\ &\quad \left. - \frac{i}{\hbar} q V \Delta t \right] \langle X_{N-1}, t_{N-1} | X_1, t_1 \rangle\end{aligned}$$

$(X_N, t_N)$  : final state

$$X_N - X_{N-1} \equiv \xi, \quad t_N - t_{N-1} \equiv \Delta t$$

then,

$$\begin{aligned}\langle X, t + \Delta t | X_0, t_0 \rangle &= \sqrt{\frac{m}{2\pi i \hbar \Delta t}} \int d\xi \exp \left( \frac{i m \xi^2}{2\hbar \Delta t} \right) \exp \left( \frac{i}{\hbar} q A \xi \right) \exp \left( -\frac{i}{\hbar} q V \Delta t \right) \\ &\quad \times \langle X - \xi, t | X, t \rangle\end{aligned}$$

since  $\xi, \Delta t$  are small,

$$\begin{aligned}&= \sqrt{\frac{m}{2\pi i \hbar \Delta t}} \int d\xi \exp \left( \frac{i m \xi^2}{2\hbar \Delta t} \right) \left[ 1 + \frac{i}{\hbar} q A \xi + \left( \frac{i}{\hbar} \right)^2 q^2 A^2 \frac{\xi^2}{2} \right] \left[ 1 - \frac{i}{\hbar} q V \Delta t \right] \\ &\quad \times \langle X - \xi, t | X, t \rangle\end{aligned}$$

$$\langle X, t + \Delta t | X, t \rangle + \Delta t \frac{\partial}{\partial t} \langle X, t | X, t \rangle = [ \dots ] \left[ \langle X, t | X, t \rangle - \xi \frac{\partial}{\partial X} \langle \dots \rangle + \frac{\xi^2}{2} \frac{\partial^2}{\partial X^2} \langle \dots \rangle \right]$$

$$\sum_{\alpha} \psi_{\alpha}^*(x, t) \left[ 1 + \Delta t \frac{\partial}{\partial t} \right] \psi_{\alpha}(x, t)$$

$$= \sum_{\alpha} \psi_{\alpha}^*(x, t) [ \dots ] \left[ 1 - \xi \frac{\partial}{\partial X} + \frac{\xi^2}{2} \frac{\partial^2}{\partial X^2} \right] \psi_{\alpha}(x, t)$$

Compare left-hand side and right-hand side,  
order  $\Delta t$

$$(l.h.s) = \Delta t \frac{\partial}{\partial t} \psi(x,t)$$

$$(r.h.s) = \sqrt{\frac{m}{2\pi i \hbar \Delta t}} \int_{-\infty}^{\infty} d\xi e^{\frac{i m \xi^2}{2 \hbar \Delta t}} \left[ \left\{ \frac{1}{2} \frac{\partial^2}{\partial x^2} - \frac{i}{\hbar} \frac{\partial}{\partial x} \phi A + \frac{1}{2} \left( \frac{i}{\hbar} \right)^2 \phi^2 A^2 \right\} \xi^2 - \frac{i q V \Delta t}{\hbar} \right] \psi(x,t)$$

NOT acting on A.

**NOTE**  $\int d\xi e^{\frac{i m \xi^2}{2 \hbar \Delta t}} = \sqrt{\frac{2\pi i \hbar \Delta t}{m}}$

$$\int d\xi \xi e^{\frac{i m \xi^2}{2 \hbar \Delta t}} = 0$$

$$\int d\xi \xi^2 e^{\frac{i m \xi^2}{2 \hbar \Delta t}} = \sqrt{\frac{2\pi i \hbar \Delta t}{m}} \left( \frac{i \hbar \Delta t}{m} \right)$$

$$(r.h.s) = \left[ \frac{1}{2} \left( \frac{i \hbar \Delta t}{m} \right) \left( \frac{\partial^2}{\partial x^2} - 2 \left( \frac{i}{\hbar} \frac{\partial}{\partial x} \right) (\phi A) + \left( \frac{i}{\hbar} \right)^2 \phi^2 A^2 \right) - \frac{i}{\hbar} V \Delta t \right] \psi(x,t)$$

$$= \left[ \frac{i \hbar}{2m} \left( \frac{\partial}{\partial x} - \frac{i}{\hbar} \phi A \right)^2 - \frac{i q V}{\hbar} \right] \Delta t \psi(x,t)$$

$$= \left[ \frac{i \hbar}{2m} \frac{i^2}{\hbar^2} \left( -i \hbar \frac{\partial}{\partial x} - \phi A \right)^2 - \frac{i q V}{\hbar} \right] \Delta t \psi(x,t)$$

$$\therefore i \hbar \partial_t \psi = \left[ \frac{1}{2m} \left( -i \hbar \frac{\partial}{\partial x} - \phi A \right)^2 + q V \right] \psi(x,t)$$

↓ 3D

$$i \hbar \partial_t \psi = \left[ \frac{1}{2m} \left( -i \hbar \vec{\nabla} - q \vec{A} \right)^2 + q V \right] \psi(x,t)$$



(ii)

$$A = \frac{B_0}{2}(-y, x, 0), \quad V = 0$$

$$\dot{z}_{cl}(0) = 0 \quad ; \quad z_{cl}(\tau) = z_{cl}(0)$$

$$\textcircled{1} S_{cl}(2,1) = \int_0^\tau dt \frac{m}{2} (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) + \frac{qB}{2} (-\dot{x}y + y\dot{x})$$

E.O.M

$$\ddot{x} - \frac{qB_0}{2m} y = 0$$

$$\ddot{y} + \frac{qB_0}{2m} x = 0$$

$$\ddot{z} = 0$$

i.e.

$$x(t) = A \cos(\omega t/2) + B \sin(\omega t/2) + C$$
$$y(t) = B \cos(\omega t/2) - A \sin(\omega t/2) + D$$

with

$$x(\tau) = x_2, \quad x(0) = x_1$$
$$y(\tau) = y_2, \quad y(0) = y_1$$
$$\omega = qB_0/m$$

$$A = -\frac{1}{2} \left( (x_2 - x_1) + \frac{\sin \omega \tau/2}{1 - \cos \omega \tau/2} (y_2 - y_1) \right)$$

$$B = \frac{1}{2} \left( -(y_2 - y_1) + \frac{\sin \omega \tau/2}{1 - \cos \omega \tau/2} (x_2 - x_1) \right)$$

$$C = \frac{1}{2} \left( (x_2 + x_1) + \frac{\sin \omega \tau/2}{(1 - \cos \omega \tau/2)} (y_2 - y_1) \right)$$

$$D = \frac{1}{2} \left( (y_2 + y_1) - \frac{\sin \omega \tau/2}{(1 - \cos \omega \tau/2)} (x_2 - x_1) \right)$$

Now, go back to the action

$$\int_0^T dt \frac{m}{2} [x^2 + y^2 - \omega xy + \omega yx]$$

$$= \frac{m \cdot \omega^2}{2} \int_0^T dt (A^2 + B^2) - (A^2 + B^2) - (BD + AC) \cos \omega t/2 + (AD - BC) \sin \omega t/2$$

$$BC = \frac{1}{4} \left[ \frac{\sin}{1 - \cos} (x_2 - x_1)(x_2 + x_1) - \frac{\sin}{1 - \cos} (y_2 - y_1)^2 - (y_2 - y_1)(x_2 + x_1) \right. \\ \left. + \left( \frac{\sin}{1 - \cos} \right)^2 (x_2 - x_1)(y_2 - x_1) \right]$$

$$AD = -\frac{1}{4} \left[ \frac{\sin}{1 - \cos} (y_2 - y_1)(y_2 + y_1) - \frac{\sin}{1 - \cos} (x_2 - x_1)^2 + (x_2 - x_1)(y_2 + y_1) \right. \\ \left. - \left( \frac{\sin}{1 - \cos} \right)^2 (x_2 - x_1)(y_2 - x_1) \right]$$

$$AD - BC = -\frac{1}{4} \left[ \frac{\sin}{1 - \cos} \left\{ (y_2 - y_1)(y_2 + y_1) - (x_2 - x_1)^2 - (y_2 - y_1)^2 \right. \right. \\ \left. \left. + (x_2 - x_1)(x_2 + x_1) \right\} - (y_2 - y_1)(x_2 + x_1) \right. \\ \left. + (x_2 - x_1)(y_2 + y_1) \right]$$

$$AC+BD = \frac{1}{4} \left[ - \left( \frac{\sin \theta}{1-\cos \theta} \right)^2 (x_2-x_1)^2 + (y_2-y_1)^2 \right]$$

$$+ \left( \frac{\sin \theta}{1-\cos \theta} \right) \left\{ (x_2-x_1)(y_2+y_1) - (y_2-y_1)(x_2+x_1) \right\}$$

$$- (y_2-y_1)(y_2+y_1) - (x_2-x_1)(x_2+x_1) \Big].$$

$$- (BD+AC)(\cos \theta/2 - 1) + (AD-BC) \sin \theta/2.$$

$$\therefore S_{ce}(2,1) = \frac{mW^2}{4} \cot \theta/2 \left[ |\vec{r}_2 - \vec{r}_1|^2 - \hat{z}(\vec{r}_2 \times \vec{r}_1) \right].$$

$$\textcircled{2} A(\tau) = \int_{x(0)=0}^{x(\tau)=0} D[x(t)] e^{\frac{i}{\hbar} \int_0^\tau dt \left[ \frac{1}{2} m \dot{x}^2 + q \dot{x} A \right]}$$

$$A_z(\tau) = \sqrt{\frac{m}{2\pi i \hbar \tau}} \quad (\text{free particle})$$

$$A_{x,y}(\tau) = \int_{\vec{x}(0)=0}^{\vec{x}(\tau)=0} D[\vec{x}(t)] e^{\frac{i}{\hbar} \int_0^\tau dt \left[ \frac{1}{2} m \dot{\vec{x}}^2 + q \dot{\vec{x}} \cdot \vec{A} \right]}, \quad \vec{x} = (x, y) \quad 2\text{-dim.}$$

$$\frac{m}{2} \dot{\vec{x}}^2 + q \dot{\vec{x}} \cdot \vec{A} = \frac{m}{2} \left[ \frac{d}{dt} \vec{x} \cdot \dot{\vec{x}} - \ddot{\vec{x}} \cdot \vec{x} \right] + \frac{q B_0}{2} [(x, y) \cdot (-y, x)]$$

↗, Boundary Condition

$$= \frac{m}{2} \left[ -\vec{x} \cdot \frac{\partial^2}{\partial t^2} \vec{x} \right]$$

$$= \frac{m}{2} \left[ -(x \ y) \begin{pmatrix} \partial_t^2 & 0 \\ 0 & \partial_t^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} + (x \ y) \begin{pmatrix} 0 & \omega \partial_t \\ -\omega \partial_t & 0 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} \right]$$

$$= \frac{m}{2} (x \ y) \underbrace{\begin{pmatrix} -\partial_t^2 & \omega \partial_t \\ -\omega \partial_t & -\partial_t^2 \end{pmatrix}}_{\equiv D} \begin{pmatrix} x \\ y \end{pmatrix}$$

$\equiv D$

$$A_{\vec{x}}(\tau) = \left( \frac{m}{2\pi i \hbar \tau} \right) \frac{1}{\det D}$$

$$= \left( \frac{1}{2\pi i \hbar \tau} \right) \prod_{n=1}^{\infty} \lambda_n, \quad \lambda_n \text{ are eigenvalues of } D.$$



$$\begin{pmatrix} -\partial_t^2 - \lambda_n & \omega \partial_t \\ -\omega \partial_t & -\partial_t^2 - \lambda_n \end{pmatrix} \begin{pmatrix} x_n \\ y_n \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$\lambda_n^+ = \partial_t^2 - \omega$$

$$= \frac{2n\pi}{\tau} - \omega$$

Because of boundary conditions  
 $\vec{x}(\tau) = \vec{x}(0)$

$$\lambda_n^- = \frac{2n\pi}{\tau} + \omega$$

$$\lambda_n \equiv \lambda_n^+ \lambda_n^-$$

$$= \left(\frac{2n\pi}{\tau}\right)^2 - \omega^2$$

$$\rightarrow 1 - \left(\frac{\omega\tau}{2n\pi}\right)^2$$

$$\prod_n \lambda_n = \prod_{n=1}^{\infty} \left[ 1 - \left(\frac{\omega\tau}{2n\pi}\right)^2 \right]$$

**NOTE**  $\frac{\sin \pi z}{\pi z} = \prod_{n=1}^{\infty} \left[ 1 - \frac{z^2}{n^2} \right]$

$$\therefore \prod_n \lambda_n = \frac{\omega\tau/2}{\sin \omega\tau/2}$$

$$A(\tau) = \left(\frac{m}{2\pi k \tau}\right)^{3/2} \frac{\omega\tau}{2 \sin \omega\tau/2}$$

#2. (i)

$$f_{\mathbf{k}}(\Omega) = -\frac{2m}{4\pi\hbar^2} \int d^d x e^{i(\mathbf{k}-\mathbf{k}') \cdot \mathbf{x}'} v(\mathbf{x}')$$

$$\frac{d\sigma}{d\Omega} = |f_{\mathbf{k}}(\Omega)|^2$$

$$\sigma = \int |f_{\mathbf{k}}(\Omega)|^2 d\Omega$$

①  $d=1$ ,  $V(r) = V_0 e^{-\kappa r}$

$$f_{\mathbf{k}} = -\frac{2m}{4\pi\hbar^2} \int_0^{\infty} dr e^{i(k-k')r} V_0 e^{-\kappa r}$$

$$= -\frac{2mV_0}{4\pi\hbar^2} \int_0^{\infty} dr e^{-(\kappa - i(k-k'))r}$$

$$= \frac{+2mV_0}{4\pi\hbar^2} \frac{e^{-\{\kappa - i(k-k')\}r}}{\kappa - i(k-k')} \Big|_0^{\infty}$$

$$= -\frac{2mV_0}{4\pi\hbar^2} \cdot \frac{1}{\kappa - i(k-k')}$$

$$|f_{\mathbf{k}}|^2 = \left| \frac{2mV_0}{4\pi\hbar^2} \right|^2 \cdot \frac{2\kappa}{\kappa^2 + (k-k')^2} = \sigma \quad \text{No Angle Dependence}$$



$$\textcircled{2} \quad d=2, \quad V(r) = \frac{V_0}{\sqrt{r}} e^{-kr}$$

$$f_{\mathbf{k}} = -\frac{2m}{4\pi\hbar^2} \int d^2\mathbf{r}' e^{i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}'} \frac{V_0}{\sqrt{r'}} e^{-kr'}$$

$$= -\frac{2mV_0}{4\pi\hbar^2} \int d^2\mathbf{r}' e^{i|\mathbf{k}-\mathbf{k}'|r'\cos\theta'} \frac{e^{-kr'}}{\sqrt{r'}}$$

$$= -\frac{2mV_0}{4\pi\hbar^2} \int_0^\infty r' dr' \int_0^{2\pi} d\theta' e^{i|\mathbf{k}-\mathbf{k}'|r'\cos\theta'} \frac{e^{-kr'}}{\sqrt{r'}}$$

~~③~~

$$= -\frac{2mV_0}{4\pi\hbar^2} F(|\mathbf{k}-\mathbf{k}'|, k)$$

$$|f_{\mathbf{k}}|^2 = \left| \frac{2mV_0}{4\pi\hbar^2} \right|^2 |F|^2$$

$$\sigma = \left| \frac{2mV_0}{4\pi\hbar^2} \right|^2 \int_0^{2\pi} d\theta |F|^2$$

$$\textcircled{3} \quad d=3, \quad V = \frac{V_0}{r} e^{-kr}$$

$$f_{\mathbf{k}}(\Omega) = - \frac{2m}{4\pi\hbar^2} \int d^3r' e^{i(\vec{k}-\vec{k}') \cdot \vec{r}'} \frac{V_0}{r'} e^{-kr'}$$

$$= - \frac{2mV_0}{4\pi\hbar^2} \cdot 2\pi \int d\cos\theta' dr' r' \cancel{r'} e^{i|\vec{k}-\vec{k}'|r'\cos\theta'} e^{-kr'}$$

$$\text{let } |\vec{k}-\vec{k}'| \equiv q$$

$$= - \frac{mV_0}{\hbar^2} \int_0^\infty dr' \frac{r'}{iqr'} \otimes 2i \sin qr' e^{-kr'}$$

$$= \frac{-2mV_0}{q\hbar^2} \int_0^\infty dr' \sin qr' e^{-kr'}$$

$$= \frac{-2mV_0}{q\hbar^2} \cdot \frac{q}{q^2+k^2}$$

$$= \frac{-2mV_0}{\hbar^2 [k^2+q^2]} \quad ; \quad q = k \sin \theta/2, \quad \theta = \text{scattering angle.}$$

$$|f_{\mathbf{k}}|^2 = \frac{4m^2V_0^2}{\hbar^4} \cdot \frac{1}{(k^2+q^2)^2}$$

$$\sigma = \frac{4m^2V_0^2}{\hbar^4} \int d\cos\theta \frac{1}{(k^2+q^2)^2}$$

(ii)

$$d=1, V = V_0 e^{-\kappa r}$$

$$\langle k' | T | k \rangle = \langle k' | V | k \rangle + \langle k' | V \frac{1}{E - H_0} V | k \rangle + \dots$$

$$\textcircled{1} \langle k' | V | k \rangle = \int dx_1 \langle k' | x_1 \rangle \langle x_1 | V | k \rangle$$

$$= \int dx_1 e^{-ik'x_1} V(x_1) e^{ikx_1}$$

$$= V_0 \int dx_1 e^{-[\kappa - i(k-k')]x_1}$$

$$= \frac{-V_0}{-i(k-k') + \kappa}$$

$$\textcircled{2} \langle k' | V \frac{1}{E - H_0} V | k \rangle$$

$$= \int dx_1 dx_2 e^{-ik'x_2} V_0 e^{-\kappa x_2} \left[ -\frac{i\mathcal{M}}{\hbar^2 k} e^{-i\mathcal{Q}k(x_1-x_2)} \right] V_0 e^{-\kappa x_1} e^{ikx_1}$$

$$= \int dx_1 dx_2 e^{-x_2(i\mathcal{M} + \kappa - ik)} V_0^2 \left( \frac{i\mathcal{M}}{\hbar^2 k} \right) e^{-\kappa x_1}$$

$$= \left( -\frac{i\mathcal{M}}{\hbar^2 k} \right) V_0^2 \left( \frac{-1}{\kappa} \right) \left( \frac{-1}{\kappa - i(k-k')} \right)$$



$$\langle k' | T | k \rangle = \frac{-V_0}{\hbar^2 k^2} \left[ 1 + \left( \frac{i m}{\hbar^2 k^2} \right) V_0 + \dots \right]$$

$$\text{if } \left| \frac{i m}{\hbar^2 k^2} V_0 \right| \ll 1,$$

$$= \frac{-V_0}{\hbar^2 k^2} \cdot \frac{1}{1 - \frac{i m V_0}{\hbar^2 k^2}}$$