## 2nd Assignment

February 2, 2017

1(a)

figure 1pt The figure should show detail of how you define your angles, and it should make sense with your initial conditions later when you solve the Lagrange's equations.

We make an ansatz here that the angles are measured from the three equi-spaced points in the figure. Then, positions of 3 masses in polar coordinates, $(r, \theta)$

$$
\left(R_{0}, \theta_{1}\right),\left(R_{0}, \theta_{2}\right),\left(R_{0}, \theta_{3}\right)
$$

The angels between masses ${ }^{1}$ are

$$
2 \pi / 3+\theta_{2}-\theta_{1}, 2 \pi / 3+\theta_{3}-\theta_{2}, 2 \pi / 3+\theta_{1}-\theta_{3}
$$

The tension of springs are acting on the circle. We can solve this system as a periodical, 3 masses, 1 dimensional coupled oscillation.

The potential energy

$$
V=\frac{1}{2} k_{0} R_{0}^{2}\left(\left(2 \pi / 3+\theta_{2}-\theta_{1}\right)^{2}+\left(2 \pi / 3+\theta_{3}-\theta_{2}\right)^{2}+\left(2 \pi / 3+\theta_{1}-\theta_{3}\right)^{2}\right)
$$

Minimize potential.

$$
\begin{gathered}
\partial_{\theta_{1}} V=k_{0} R_{0}^{2}\left(2 \theta_{1}-\theta_{2}-\theta_{3}\right)=0 \\
\partial_{\theta_{2}} V=k_{0} R_{0}^{2}\left(-\theta_{1}-\theta_{2}+2 \theta_{3}\right)=0 \\
\partial_{\theta_{3}} V=k_{0} R_{0}^{2}\left(-\theta_{1}-\theta_{2}+2 \theta_{3}\right)=0
\end{gathered}
$$

Therefore, $\theta_{1}=\theta_{2}=\theta_{3}=2 \pi / 3$. Our ansatz makes sense. 1 pt for the proper reasoning and result

[^0]The kinetic energy is

$$
T=\frac{m R_{0}^{2}}{2}\left(\dot{\theta}_{1}^{2}+\dot{\theta}_{2}^{2}+\dot{\theta}_{3}^{2}\right)
$$

The Lagrange's equations

$$
\begin{gathered}
m R_{0}^{2} \ddot{\theta}_{1}+k_{0} R_{0}^{2}\left(2 \theta_{1}-\theta_{2}-\theta_{3}\right)=0 \\
m R_{0}^{2} \ddot{\theta}_{2}+k_{0} R_{0}^{2}\left(-\theta_{1}+2 \theta_{2}-\theta_{3}\right)=0 \\
m R_{0}^{2} \ddot{\theta}_{3}+k_{0} R_{0}^{2}\left(-\theta_{1}-\theta_{2}+2 \theta_{3}\right)=0
\end{gathered}
$$

2 (b)
Let $\theta_{i}=A_{i} \exp (i \omega t)$

$$
\begin{gathered}
-\omega^{2} m R_{0}^{2} A_{1}+k_{0} R_{0}^{2}\left(2 A_{1}-A_{2}-A_{3}\right)=0 \\
-\omega^{2} m R_{0}^{2} A_{2}+k_{0} R_{0}^{2}\left(-A_{1}+2 A_{2}-A_{3}\right)=0 \\
-\omega^{2} m R_{0}^{2} A_{3}+k_{0} R_{0}^{2}\left(-A_{1}-A_{2}+2 A_{3}\right)=0 \\
M=\left(\begin{array}{ccc}
2 k_{0} R_{0}^{2}-m R_{0}^{2} \omega^{2} & -k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} \\
-k_{0} R_{0}^{2} & 2 k_{0} R_{0}^{2}-m R_{0}^{2} \omega^{2} & -k_{0} R_{0}^{2} \\
-k_{0} R_{0}^{2} & k_{0} R_{0}^{2} & 2 k_{0} R_{0}^{2}-m R_{0}^{2} \omega^{2}
\end{array}\right) \\
\operatorname{Det}(M)=m R_{0}^{6} \omega^{2}\left(-m^{2} \omega^{4}+6 k_{0} m \omega^{2}-9 k_{0}^{2}\right) \\
\omega=0, \pm \sqrt{\frac{3 k_{0}}{m}} 1 \mathrm{pt}
\end{gathered}
$$

For $\omega=0$,

$$
M=\left(\begin{array}{ccc}
2 k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} \\
-k_{0} R_{0}^{2} & 2 k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} \\
-k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} & 2 k_{0} R_{0}^{2}
\end{array}\right)
$$

$\omega=0$ leads linear motion, so the corresponding solution is

$$
\left(c_{1}+c_{2} t\right)\left(\begin{array}{l}
A_{1} \\
A_{2} \\
A_{3}
\end{array}\right)=\left(c_{1}+c_{2} t\right)\left(\begin{array}{l}
1 \\
1 \\
1
\end{array}\right)
$$

For $\omega= \pm \sqrt{\frac{3 k_{0}}{m}}$,

$$
M=\left(\begin{array}{ccc}
-k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} \\
-k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} \\
-k_{0} R_{0}^{2} & -k_{0} R_{0}^{2} & -k_{0} R_{0}^{2}
\end{array}\right)
$$

Here, we have a condition, $A_{1}+A_{2}+A_{3}=0$, and two corresponding orthogonal eigenvectors. Without loss of generality, $A_{1}=1$. We can choose two $A_{2}$ and $A_{3}$ which satisfy $A_{2}+A_{3}=-1$. From intuition of the initial condition we will meet later and some experience of 1- dimensional 3 masses coupled oscillations, we can choose $A_{2}=0$, and, therefore, $A_{3}=-1$. As the result of the choice, the remained normal mode must be $(1,-2,1)$ because it should be orthogonal to $(1,1,1)$ and $(1,0,-1)$. Because we have $\omega= \pm \sqrt{\frac{3 k_{0}}{m}}$, the resultant oscillation should be the combination of $\exp \left(\sqrt{\frac{3 k_{0}}{m}}\right)$ and $\exp \left(-\sqrt{\frac{3 k_{0}}{m}}\right)$. This can be replaced by the combination of sine and cosine.

Thus, corresponding and orthogonal solutions are

$$
\left(c_{3} \cos \sqrt{\frac{3 k_{0}}{m}} t+c_{4} \sin \sqrt{\frac{3 k_{0}}{m}} t\right)\left(\begin{array}{c}
1 \\
0 \\
-1
\end{array}\right)
$$

and

$$
\left(c_{5} \cos \sqrt{\frac{3 k_{0}}{m}} t+c_{6} \sin \sqrt{\frac{3 k_{0}}{m}} t\right)\left(\begin{array}{c}
1 \\
-2 \\
1
\end{array}\right)
$$

The full solution is ( 1 pt for proper orthogonal normal modes)
$\left(\begin{array}{l}\theta_{1} \\ \theta_{2} \\ \theta_{3}\end{array}\right)=\left(c_{1}+c_{2} t\right)\left(\begin{array}{l}1 \\ 1 \\ 1\end{array}\right)+\left(c_{3} \cos \sqrt{\frac{3 k_{0}}{m}} t+c_{4} \sin \sqrt{\frac{3 k_{0}}{m}} t\right)\left(\begin{array}{c}1 \\ 0 \\ -1\end{array}\right)+\left(c_{5} \cos \sqrt{\frac{3 k_{0}}{m}} t+c_{6} \sin \sqrt{\frac{3 k_{0}}{m}} t\right)\left(\begin{array}{c}1 \\ -2 \\ 1\end{array}\right)$
(i) 1 pt for both initial conditions

Stationary condition yields the time derivative initial condition

$$
\dot{\dot{\theta}_{0}}=\left(\begin{array}{c}
\dot{\theta}_{01} \\
\dot{\theta}_{02} \\
\dot{\theta}_{03}
\end{array}\right)=\left(\begin{array}{l}
0 \\
0 \\
0
\end{array}\right)
$$

(ii)

From the above figure, we have set three equilibrium points are at $(0,2 \pi / 3,4 \pi / 3) . m_{1}$ is at 0 , and both $m_{2}$ and $m_{3}$ are at the opposite position on a circle. I have a positive sign along clockwise direction. By my convention, the angle of $m_{2}$ is $+\pi / 3$ from the 2 nd equilibrium point. Similarly, $m_{3}$ placed at $+2 \pi / 3$ from the 3 rd equilibrium point.

$$
\vec{\theta}_{0}=\left(\begin{array}{c}
\theta_{01} \\
\theta_{02} \\
\theta_{03}
\end{array}\right)=\left(\begin{array}{c}
0 \\
\frac{\pi}{3} \\
\frac{2 \pi}{3}
\end{array}\right)
$$

By time derivative initial condition, $c_{2}=c_{4}=c_{6}=0$, by the angle initial condition, $c_{1}=\pi / 3$, $c_{3}=c_{5}=-\pi / 6$

The final solution is (1pt)

$$
\left(\begin{array}{c}
\theta_{1} \\
\theta_{2} \\
\theta_{3}
\end{array}\right)=\left(\begin{array}{c}
\pi / 3\left(1-\cos \sqrt{\frac{3 k_{0}}{m}} t\right) \\
\pi / 3 \\
\pi / 3\left(1+\cos \sqrt{\frac{3 k_{0}}{m}} t\right)
\end{array}\right)
$$

If your solution contains cosine and minus cosine with the correct coefficients, you get full marks. For example, if you use $(\pi / 3,0,-\pi / 3)$ as an initial condition, you will get $\vec{\theta}=\pi / 3 \cos \sqrt{\frac{3 k_{0}}{m}} t(-1,0,1)$. Both solutions physically indicate that one mass is at rest and other two oscillate in opposite direction with same magnitude of the amplitude. These are sum of two normal modes, $(1,1,1)$ and ( $1,0,-1$ ).
2.
(i) mass : kg , spring constant $: \mathrm{kg} / \mathrm{s}^{2}$ (1pt)
(a) (1pt)
$z_{1}: \vec{F}=-k \vec{z}_{1}=-2 \times 5 \times 10^{3} \mathrm{~m} \cdot \mathrm{~kg} / \mathrm{s}^{2} \hat{z}=-10^{4} \mathrm{~m} \cdot \mathrm{~kg} / \mathrm{s}^{2} \hat{z}$
$z_{2}: \vec{F}=-k \vec{z}_{2}=-8 \times 5 \times 10^{3} \mathrm{~m} \cdot \mathrm{~kg} / \mathrm{s}^{2} \hat{z}=-4 \times 10^{4} \mathrm{~m} \cdot \mathrm{~kg} / \mathrm{s}^{2} \hat{z}$
(b)
$z_{1}: E=\frac{1}{2} k z_{1}^{2}=\frac{1}{2} \times 5 \times 10^{3} \times 4 m^{2}=10^{4} J$
$z_{2}: E=\frac{1}{2} k z_{1}^{2}=\frac{1}{2} \times 5 \times 10^{3} \times 64 m^{2}=1.6 \times 10^{5} \mathrm{~J}$
$\omega=\sqrt{\frac{k}{M}}=\sqrt{\frac{5 \times 10^{3} \mathrm{~kg} / \mathrm{s}^{2}}{2 \times 10^{4} \mathrm{~kg}}}=1 / 2 \mathrm{~Hz}$ ( 1 pt for energy and angular frequency)
$\Delta z(t)=A \cos \omega t$
Kinetic Energy : $K=\frac{1}{2} M \Delta \dot{z}^{2}=\frac{1}{2} M A^{2} \omega^{2} \sin ^{2} \omega t$
Potential energy by gravity is canceled by floating.
Power: $P=\frac{d K}{d t}=\frac{1}{2} M \Delta \dot{z}^{2}=\frac{1}{2} M A^{2} \omega^{3} \sin 2 \omega t$
Period : $T=2 \pi / \omega$
Mean Power : ( 1 pt for power and mean power )
Because the all kinetic energy are extracted.
$\frac{1}{T} \int_{0}^{T}|P| d t=\frac{1}{2} M A^{2} \omega^{3} \times \frac{2}{\pi}=1 / 2 \times 2 \times 10^{4} \mathrm{~kg} \times 2^{2} \mathrm{~m}^{2} \times 1 / 2^{3} \mathrm{H} z^{3} \times \frac{2}{\pi}=3183.1$ Watt. ${ }^{2}$
Besides, mean power of sine function is obtained from the root mean square. I will give you full marks for this, too.

$$
\sqrt{\frac{1}{T} \int_{0}^{T} P^{2} d t}=\frac{1}{2} M A^{2} \omega^{3} \times \sqrt{\frac{1}{2}}=1 / 2 \times 2 \times 10^{4} k g \times 2^{2} m^{2} \times 1 / 2^{3} H z^{3} \times \sqrt{\frac{1}{2}}=3535.53
$$

Watt. ${ }^{3}$

[^1]
[^0]:    ${ }^{1}$ The relative angles from the 3 spaced points

[^1]:    ${ }^{2}$ Mean of $|\sin t|$ is $\sqrt{\frac{2}{\pi}}$. In other words, $\frac{1}{2 \pi} \int_{0}^{2 \pi}|\sin t| d t=\sqrt{\frac{2}{\pi}}$
    ${ }^{3}$ RMS of $\sin t$ is $\sqrt{\frac{1}{2}}$.

