

PROBLEM SET | SOLUTIONS

We can use the same change of variables as on the worksheet,

$$\phi(x,t) = \sum_n \phi_n(t) \sin\left(\frac{n\pi x}{L}\right).$$

The equations of motion become:

$$\rho \sum_n \ddot{\phi}_n \sin\left(\frac{n\pi x}{L}\right) = \tau \sum_n \phi_n(t) \left\{ -\left(\frac{n\pi}{L}\right)^2 \sin\left(\frac{n\pi x}{L}\right) \right\} \\ - \mu \sum_n \phi_n(t) \left\{ \sin\left(\frac{n\pi x}{L}\right) \right\}$$

$$\Rightarrow \sum_n \sin\left(\frac{n\pi x}{L}\right) \left\{ \rho \ddot{\phi}_n + \left\{ \tau \left(\frac{n\pi}{L}\right)^2 + \mu \right\} \phi_n \right\} = 0$$

The Fourier decomposition of ϕ has all terms equal to 0, so we find

$$\ddot{\phi}_n = -\frac{1}{\rho} \left\{ \tau \left(\frac{n\pi}{L}\right)^2 + \mu \right\} \phi_n.$$

The variable ϕ_n has the same e.o.m. as the displacement of a harmonic oscillator with freq. $\omega_n = \sqrt{\frac{1}{\rho} \left(\tau \left(\frac{n\pi}{L}\right)^2 + \mu \right)}$.

For the energy, the change of variables yields

$$E = \int_0^L dx \left\{ \frac{1}{2} \rho \left(\sum_n \dot{\phi}_n \sin\left(\frac{n\pi x}{L}\right) \right) \left(\sum_m \dot{\phi}_m \sin\left(\frac{m\pi x}{L}\right) \right) \right. \\ \left. + \frac{1}{2} \tau \left(\sum_n \phi_n \left(\frac{n\pi}{L}\right) \cos\left(\frac{n\pi x}{L}\right) \right) \left(\sum_m \phi_m \left(\frac{m\pi}{L}\right) \cos\left(\frac{m\pi x}{L}\right) \right) \right. \\ \left. + \frac{1}{2} \mu \left(\sum_n \phi_n \sin\left(\frac{n\pi x}{L}\right) \right) \left(\sum_m \phi_m \sin\left(\frac{m\pi x}{L}\right) \right) \right\}$$

$$= \sum_n \frac{\rho L}{4} \dot{\phi}_n^2 + \left\{ \frac{\tau L}{4} \left(\frac{n\pi}{L} \right)^2 + \frac{\mu L}{4} \right\} \phi_n^2$$

where we have used

$$\int_0^L dx \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi x}{L}\right) = \frac{L}{\pi} \int_0^\pi dy \sin(ny) \sin(my) \\ = \frac{L}{2} \times S_{n,m}$$

$$\text{and } \int_0^L dx \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi x}{L}\right) = \frac{L}{2} S_{n,m}.$$

This matches with the energy for a bunch of harmonic oscillators

$$E = \sum_n \frac{1}{2} m_n \dot{x}_n^2 + \frac{1}{2} m_n \omega_n^2 x_n^2$$

if we define $m_n = \frac{\rho L}{2}$

$$\omega_n = \sqrt{\frac{\tau}{\rho} \left(\frac{n\pi}{L} \right)^2 + \frac{\mu}{\rho}}.$$

The field theory is mathematically equivalent to a system of harmonic oscillators at the classical level, so the quantum spectrum for the two systems is the same. For the subsystem described by ϕ_n , the allowed quantum energies are:

$$\begin{aligned} E - E_0 &= \hbar\omega_n N_n \\ &= \hbar N_n \sqrt{\frac{c}{\rho} \left(\frac{n\pi}{L}\right)^2 + \frac{\mu}{\rho}} \end{aligned}$$

The allowed energies for the full system are

$$E - E_0 = \sum_n \hbar N_n \sqrt{\frac{c}{\rho} \left(\frac{n\pi}{L}\right)^2 + \frac{\mu}{\rho}}$$

- b) For particles of mass M , the energies are given by

$$E = \sqrt{M^2 c^4 + p^2 c^2}$$

the allowed momenta are such that the de Broglie wavelength $\lambda = \frac{h}{p}$ is $2L, L, \frac{2L}{3}, \frac{L}{2}, \dots$ (we want an integer or half-integer number of wavelengths to fit in the box. Thus, the allowed energies for a single particle are

$$E_n = \sqrt{M^2 c^4 + c^2 \cdot \frac{h^2}{\lambda_n^2}} = \sqrt{M^2 c^4 + \frac{c^2 h^2 n^2}{(2L)^2}}$$

We can have any number of particles at each wave length (* if they are bosons *) so the total energy can be :

$$E = \sum_n N_n \cdot \sqrt{M^2 c^4 + \frac{c^2 h^2 n^2}{(2L)^2}} \quad (\text{A})$$

4) If we take $\frac{\epsilon}{p} = c^2$, the field theory energies are :

$$E - E_0 = \sum_n N_n \sqrt{\frac{\hbar^2 \mu}{p} + \frac{c^2 h^2 n^2}{(2L)^2}}$$

This matches exactly with (A) if we take

$$\mu = \frac{p M^2 c^4}{\hbar^2}$$